When electrons are confined to length scales approaching the de Broglie wavelength, their behavior is dominated by quantum mechanical effects. Here we report the construction and characterization of structures for confining electrons to this length scale. The walls of these "quantum corrals" are built from Fe atoms which are individually positioned on the Cu (111) surface by means of a scanning tunneling microscope (STM). These atomic structures confine surface state electrons laterally because of the strong scattering that occurs between surface state electrons and the Fe atoms. The surface state electrons are confined in the direction perpendicular to the surface because of intrinsic energetic barriers that exist in that direction.

This is the first paragraph of "Confinement of Electrons to Quantum Corrals on a Metal Surface," published by M. F. Crommie, C. P. Lutz, and D. M. Eigler in the October 8, 1993 issue of Science Magazine. They report the corraling of the surface electrons of Cu in a ring of radius 135 a₀ created by 48 Fe atoms. The quantum mechanics for this form of electron confinement is well-known. Schroedinger's equation for a particle in a ring and its solution (in atomic units) are given below.

\[
\frac{1}{2 \mu} \frac{d^2}{dr^2} \Psi(r) - \frac{1}{2 r \mu} \frac{d}{dr} \Psi(r) + \left( \frac{L^2}{2 \mu r^2} \right) \Psi(r) = E \Psi(r) \label{1}
\]

with energies

\[
|E_{n,L} = \frac{Z_{n,L}^2}{2\mu R^2} \label{2} \]

and the unnormalized wavefunctions

\[
|\Psi_{n,L} = J_z (Z_{n,L} , R) \label{3} |
\]

The calculated results are displayed by plotting the wave function squared in Cartesian coordinates. The exponential term involving L and \(\langle \theta \rangle\) is discarded because \(\langle |\mathrm{e}^{\mathrm{i} \cdot \mathrm{L} \cdot \theta}|\rangle^2 = 1\).
The theoretical results are displayed by plotting the wave function in Cartesian coordinates.

$$\begin{align*}
\mathcal{R} &= 135 \\
\mathcal{n} &= 5 \\
\mathcal{L} &= 0 \\
\mathcal{N} &= 100 \\
\mathcal{i} &= 0 \ldots \mathcal{N} \\
\mathcal{j} &= 0 \ldots \mathcal{N} \\
\mathbf{x}_i &= -\mathcal{R} + \frac{2 \cdot i}{\mathcal{N}} \cdot \mathcal{R} \\
\mathbf{y}_j &= -\mathcal{R} + \frac{2 \cdot j}{\mathcal{N}} \cdot \mathcal{R} \\
\Psi(x,y) &= \begin{cases} 
\text{if } \sqrt{x^2+y^2} \leq \mathcal{R} 
\mathcal{J}_{\mathcal{n}}(\mathcal{L}, \frac{\sqrt{x^2+y^2}}{\mathcal{R}}) \\
\text{otherwise} 
\end{cases}
\end{align*}$$

$$\quad P_{i,j} = \Psi(x_i, y_j)^2$$

Figure 1: (left) The $|5,0\rangle$ wavefunction of an electron in the 2D quantum corral. (Right) The experimental surface electron density reported by Crombie, et al. is shown below. The agreement between theory and experiment is very good.

However, Crommie, et al. noted that the $|5,0\rangle$, $|4,2\rangle$ and $|2,7\rangle$ states are close in energy, being proportional to the squares of 14.931, 14.796 and 14.81 given in the table above. An even statistical mixture of these states would yield the surface electron density shown below, which is also visually in agreement with the experimental surface electron density.

$$\begin{align*}
\Psi'(x,y) &= \begin{cases} 
\mathcal{J}_{\mathcal{n}}(0, \frac{\sqrt{x^2+y^2}}{\mathcal{R}})^2 + \mathcal{J}_{\mathcal{n}}(2, \frac{\sqrt{x^2+y^2}}{\mathcal{R}})^2 + \mathcal{J}_{\mathcal{n}}(7, \frac{\sqrt{x^2+y^2}}{\mathcal{R}})^2 
\text{if } \sqrt{x^2+y^2} \leq \mathcal{R} \\
\text{otherwise} 
\end{cases}
\end{align*}$$

$$\quad P_{i,j} := \Psi'(x_i, y_j)$$
References

1. M. F. Crommie, C. P. Lutz, and D. M. Eigler in the October 8, 1993 issue of Science Magazine

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