Q10.8

Show that the four $|sp^3\rangle$ orbitals are orthonormal.

S10.8

This means showing that each pair of $|sp^3 \rangle$ hybrid orbitals meets the criteria: $\langle sp^3_i | sp^3_j \rangle = \delta_{ij}$

Designating the four $sp^3$ orbitals as:

\[
|1\rangle = \frac{1}{2}( |s\rangle + |p_x\rangle + |p_y\rangle + |p_z\rangle )
\]

\[
|2\rangle = \frac{1}{2}( |s\rangle - |p_x\rangle - |p_y\rangle + |p_z\rangle )
\]

\[
|3\rangle = \frac{1}{2}( |s\rangle + |p_x\rangle - |p_y\rangle - |p_z\rangle )
\]

\[
|4\rangle = \frac{1}{2}( |s\rangle - |p_x\rangle + |p_y\rangle - |p_z\rangle )
\]

Normality:

\[
\langle 1 | 1 \rangle = \frac{1}{4} \left( \langle s|s \rangle + \langle s| p_x \rangle + \langle s| p_y \rangle + \langle s| p_z \rangle + \langle p_x| s \rangle + \langle p_x| p_x \rangle + \langle p_x| p_y \rangle + \langle p_x | p_z \rangle + \langle p_y| s \rangle + \langle p_y | p_x \rangle + \langle p_y | p_y \rangle + \langle p_y| p_z \rangle + \langle p_z| s \rangle + \langle p_z | p_x \rangle + \langle p_z | p_y \rangle + \langle p_z| p_z \rangle \right) = 1
\]

\[
\langle 2 | 2 \rangle = \frac{1}{4} \left( \langle s|s \rangle - \langle s| p_x \rangle - \langle s| p_y \rangle + \langle s| p_z \rangle - \langle p_x| s \rangle + \langle p_x| p_x \rangle + \langle p_x| p_y \rangle - \langle p_x | p_z \rangle - \langle p_y| s \rangle + \langle p_y | p_x \rangle + \langle p_y | p_y \rangle - \langle p_y| p_z \rangle + \langle p_z| s \rangle - \langle p_z | p_x \rangle - \langle p_z | p_y \rangle + \langle p_z| p_z \rangle \right) = 1
\]
\[ \langle 1 | 2 \rangle = \langle 2 | 1 \rangle = \left( \dfrac{1}{2}( \langle s| - \langle p_x | + \langle p_y| - \langle p_z|) \right) \left( \dfrac{1}{2}( |s \rangle + | p_x \rangle + |p_y \rangle + | p_z \rangle) \right) \]
\[= \dfrac{1}{4} \langle s|s \rangle - \langle p_x| p_x \rangle - \langle p_y | p_y \rangle + \langle p_z| p_z \rangle = \dfrac{1}{4}( 1-1-1+1) = 0 \]

Orthogonality:

For the longer wavefunctions, only surviving terms in the dot product are included:

\[ \langle 1 | 2 \rangle = \langle 2 | 1 \rangle = \left( \dfrac{1}{2}( \langle s| - \langle p_x | - \langle p_y| + \langle p_z|) \right) \left( \dfrac{1}{2}( |s \rangle + | p_x \rangle + |p_y \rangle + | p_z \rangle) \right) \]
\[= \dfrac{1}{4} \langle s|s \rangle + \langle p_x| p_x \rangle - \langle p_y | p_y \rangle - \langle p_z| p_z \rangle = \dfrac{1}{4}( 1+1-1-1) = 0 \]

\[ \langle 1 | 4 \rangle = \langle 4 | 1 \rangle = \left( \dfrac{1}{2}( \langle s| - \langle p_x | + \langle p_y| - \langle p_z|) \right) \left( \dfrac{1}{2}( |s \rangle + | p_x \rangle + |p_y \rangle + | p_z \rangle) \right) \]
\[= \dfrac{1}{4} \langle s|s \rangle - \langle p_x| p_x \rangle + \langle p_y | p_y \rangle - \langle p_z| p_z \rangle = \dfrac{1}{4}( 1-1+1-1) = 0 \]

\[ \langle 2 | 3 \rangle = \langle 3 | 2 \rangle = \left( \dfrac{1}{2}( \langle s| + \langle p_x | - \langle p_y| - \langle p_z|) \right) \left( \dfrac{1}{2}( |s \rangle - | p_x \rangle - |p_y \rangle + | p_z \rangle) \right) \]
\[= \dfrac{1}{4} \langle s|s \rangle + \langle p_x| p_x \rangle - \langle p_y | p_y \rangle - \langle p_z| p_z \rangle = \dfrac{1}{4}( 1+1+1-1) = 0 \]

Orthogonality:

For the longer wavefunctions, only surviving terms in the dot product are included:
\[ \langle 1 | 4 \rangle = \langle 4 | 1 \rangle = \left( \dfrac{1}{4}(\langle s | s \rangle - \langle p_x | p_x \rangle + \langle p_y | p_y \rangle - \langle p_z | p_z \rangle) \right) \]

\[ \langle 2 | 4 \rangle = \langle 4 | 2 \rangle = \left( \dfrac{1}{2}(\langle s | - \langle p_x | + \langle p_y | - \langle p_z |) \right) \]

\[ \langle 3 | 4 \rangle = \langle 4 | 3 \rangle = (\dfrac{1}{2}(\langle s | - \langle p_x | - \langle p_y | - \langle p_z |)) \]

\[ \langle 4 | 4 \rangle = \langle 4 | 4 \rangle = \left( \dfrac{1}{4}(\langle s | s \rangle + \langle p_x | p_x \rangle - \langle p_y | p_y \rangle - \langle p_z | p_z \rangle) \right) \]

\[ \langle 5 | 4 \rangle = \langle 4 | 5 \rangle = \left( \dfrac{1}{2}(\langle s | + \langle p_y | + \langle p_z | - \langle p_w |) \right) \]

\[ \langle 6 | 4 \rangle = \langle 4 | 6 \rangle = \left( \dfrac{1}{4}(\langle s | s \rangle + \langle p_y | p_y \rangle - \langle p_z | p_z \rangle + \langle p_w | p_w \rangle) \right) \]

Q10.9

The \(sp^{3}d^{2}\) hybrid orbitals are given by:

\[ \chi_1 (r) = \dfrac{1}{\sqrt{6}}\psi_{3s}(r) - \dfrac{1}{\sqrt{2}}\psi_{3p_x}(r) - \dfrac{1}{\sqrt{12}}\psi_{3d_{z^2}} + \dfrac{1}{\sqrt{4}}\psi_{3d_{x^2-y^2}} \]

\[ \chi_2 (r) = \dfrac{1}{\sqrt{6}}\psi_{3s}(r) + \dfrac{1}{\sqrt{2}}\psi_{3p_x}(r) - \dfrac{1}{\sqrt{12}}\psi_{3d_{z^2}} + \dfrac{1}{\sqrt{4}}\psi_{3d_{x^2-y^2}} \]

\[ \chi_3 (r) = \dfrac{1}{\sqrt{6}}\psi_{3s}(r) - \dfrac{1}{\sqrt{2}}\psi_{3p_y}(r) - \dfrac{1}{\sqrt{12}}\psi_{3d_{z^2}} - \dfrac{1}{\sqrt{4}}\psi_{3d_{x^2-y^2}} \]

\[ \chi_4 (r) = \dfrac{1}{\sqrt{6}}\psi_{3s}(r) + \dfrac{1}{\sqrt{2}}\psi_{3p_y}(r) - \dfrac{1}{\sqrt{12}}\psi_{3d_{z^2}} - \dfrac{1}{\sqrt{4}}\psi_{3d_{x^2-y^2}} \]

\[ \chi_5 (r) = \dfrac{1}{\sqrt{6}}\psi_{3s}(r) - \dfrac{1}{\sqrt{2}}\psi_{3p_z}(r) + \dfrac{1}{\sqrt{12}}\psi_{3d_{z^2}} \]

\[ \chi_6 (r) = \dfrac{1}{\sqrt{6}}\psi_{3s}(r) + \dfrac{1}{\sqrt{2}}\psi_{3p_z}(r) + \dfrac{1}{\sqrt{12}}\psi_{3d_{z^2}} \]

Determine the angles of \(SF_{6}\) using the vector approach (dot product formula).

S10.9

The \(s\) orbitals are spherical and therefore do not contribute to the directional vectors for this problem. The \(d_{z^2}\) orbital only has \(z\) directionality and the \(d_{x^2-y^2}\) orbital has equal parts \(x\) and \(y\) directionality. The equations can be rewritten:
\[ \psi_3 = -\dfrac{1}{\sqrt{4}} \mathbf{i} + \left( -\dfrac{1}{\sqrt{2}} - \dfrac{1}{\sqrt{4}} \right) \mathbf{j} - \dfrac{1}{\sqrt{12}} \mathbf{k} \]
\[ \psi_4 = -\dfrac{1}{\sqrt{4}} \mathbf{i} + \left( \dfrac{1}{\sqrt{2}} - \dfrac{1}{\sqrt{4}} \right) \mathbf{j} - \dfrac{1}{\sqrt{12}} \mathbf{k} \]
\[ \psi_5 = \left( -\dfrac{1}{\sqrt{2}} + \dfrac{1}{\sqrt{12}} \right) \mathbf{k} \]
\[ \psi_6 = \left( \dfrac{1}{\sqrt{2}} + \dfrac{1}{\sqrt{12}} \right) \mathbf{k} \]

For any two \( \psi \) the angle can be calculated using:
\[ \left( \sqrt{A_{x}^2 + A_{y}^2 + A_{z}^2} \right) \left( \sqrt{B_{x}^2 + B_{y}^2 + B_{z}^2} \right) \cos \theta = A_{x}B_{x} + A_{y}B_{y} + A_{z}B_{z} \]

So, choosing \( \psi_5 \) and \( \psi_6 \) we get:
\[ \left( -\dfrac{1}{\sqrt{2}} + \dfrac{1}{\sqrt{12}} \right) \left( \dfrac{1}{\sqrt{2}} + \dfrac{1}{\sqrt{12}} \right) \cos \theta = -\dfrac{5}{12} \]
\[ \cos \theta = 0 \]
\[ \theta = 90^\circ \]

Q10.10

Given \( \xi_1 = \dfrac{1}{\sqrt{4}}2s + \sqrt{\dfrac{3}{4}}2p_z \)
\[ \xi_2 = \dfrac{1}{\sqrt{4}}2s + \sqrt{\dfrac{3}{4}}2p_x - \dfrac{1}{\sqrt{12}}2p_z \]

what is the angle between \( \xi_1 \) and \( \xi_2 \)?

What is the purpose of the square root constants before the orbitals?

S10.10

Since we know that
\[ \left( \sqrt{A_x^2 + A_y^2 + A_z^2} \right) \left( \sqrt{B_x^2 + B_y^2 + B_z^2} \right) \cos \theta = A_xB_x + A_yB_y + A_zB_z \]

When we plug in our constants in the x, y, and z directions for the two equations we get
\[ \left( \sqrt{\dfrac{1}{2} + \dfrac{1}{12}} \right) \left( \sqrt{\dfrac{3}{4}} \right) \cos \theta = -\sqrt{\dfrac{1}{12}} \sqrt{\dfrac{3}{4}} \]

isolating the theta term we get
\[ \theta = \arccos(-\sqrt{\dfrac{1}{12}}/\sqrt{\dfrac{12}{9}}) \]


yielding

\[ \theta = 109.47^\circ \]

The square root constants are the necessary normalization constants.

Q10.11

Given that \[ \psi_1 = 0.71j + 0.55k \] and \[ \psi_2 = -0.71j + 0.55k \] Find the bond angle between \( \psi_1 \) and \( \psi_2 \) using the vector approach.

S10.11

Using this equation below:

\[
\left(\sqrt{A_x^2 + A_y^2 + A_z^2}\right) \left(\sqrt{B_x^2 + B_y^2 + B_z^2}\right) \cos \theta = A_x B_x + A_y B_y + A_z B_z
\]

use to find the angle between the orbitals,

\[
(0.71^2 + 0.55^2)^{0.5} (0.71^2 + 0.55^2)^{0.5} \cos \theta = -0.71^2 + 0.55^2
\]

\[ \cos \theta = -0.25 \]

\[ \theta = 104.5^\circ \]

Q10.12

Assuming a water molecule sits in the yz-plane, show that two bonding hybrid atomic orbitals on the oxygen atom can be expressed as

\[ \psi_1 = N[\gamma 2s + (\sin \theta) 2p_y + (\cos \theta) 2p_z]\]

and

\[ \psi_2 = N[\gamma 2s - (\sin \theta) 2p_y + (\cos \theta) 2p_z]\]

Additionally, find \( \gamma \) assuming the bond angles to be \( 104.5^\circ \).

S10.12

Because the molecules are in the yz plane, the \( p_x \) orbital can be neglected because it is completely orthogonal. The hybrid orbitals are a linear combination of the \( 2s \), \( 2p_y \), and \( 2p_z \) orbitals.

\[ \psi_{\text{bonding}} = N[\gamma 2s + c_1 2p_y + c_2 2p_z] \]
\(c_1\) and \(c_2\) can be found by thinking about how the bonds are oriented. The two bonds both have \(\cos[\theta]\) character in the z-direction and have \(\pm \sin[\theta]\) in the y-direction. The two functions can therefore be written as

\[
\psi_1 = N[\gamma 2s + (\sin[\theta]2p_y + \cos[\theta]2p_z)]
\]

and

\[
\psi_2 = N[\gamma 2s - (\sin[\theta]2p_y + \cos[\theta]2p_z)]
\]

when \(c_1\) and \(c_2\) are replaced with the sin and cos functions above.

The two functions when integrated will equal zero due to orthogonality.

\[
\langle\psi_1|\psi_2\rangle = 0
\]

\[
N^2(\gamma^2 + \cos^2[\theta] - \sin^2[\theta])
\]

\[
\gamma^2 = \sin^2[\theta] - \cos^2[\theta]
\]

The bond angles in water are 104.5° which can be substituted into \(\theta\).

\[
\gamma = 0.5
\]

---

**Q10.13**

The lone pair wave functions of H2O can be described as:

\[
\psi_{l1} = 0.54(2s) - 0.44(2pz) + 0.72(2px)
\]

\[
\psi_{l2} = 0.54(2s) - 0.44(2pz) - 0.72(2px)
\]

Confirm the orthonormality of these wave functions.

---

**S10.13**

The two wave functions must be normalized and orthogonal to each other.

\[
\int d\tau \psi_1^\ast(\psi_1) = 1
\]

\[
(0.54)^2(2s) + (-0.44)^2(2pz) + (0.72)^2(2px) = 1
\]

\[
\int d\tau \psi_2^\ast(\psi_2) = 1
\]

\[
(0.54)^2(2s) + (-0.44)^2(2pz) + (-0.72)^2(2px) = 1
\]

\[
\int d\tau \psi_1^\ast(\psi_2) = 0
\]

\[
(0.54)^2(2s) + (-0.44)^2(2pz) - (0.72)^2(2px) = 0
\]
The atomic orbitals within the linear combination of the lone pair wave functions are normalized and orthogonal, zeroing out cross products and leaving only the squares of the coefficients.

10.14

Molecular orbitals for a linear $XY_2$ molecule can be represented as

Draw a schematic representation for the $3\sigma_g$, $4\sigma_g$, $1\pi_g$, $2\sigma_g$ orbitals

Which has the highest energy?

S10.14

[$3\sigma_g$]

NOTE:

[$3\sigma_g$]

$4\sigma_g$
The $4\sigma_g$ molecular orbital has the highest energy; as expected since it more nodes.

Q10.18

Use the given Walsh diagram to predict the geometry of the following molecules:

a. $\text{H}_2\text{O}$
b. $\text{H}_2\text{S}$
c. $\text{H}_2\text{Be}$
The Walsh Diagram predicts the geometry of a molecule by assigning its valence electrons to the appropriate energy levels. In general, the lowest energy configuration is preferred.

a. \(\text{H}_2\text{O}\) has 8 valence electrons, which corresponds to the 4\(^{th}\) highest orbital on the diagram. The \textit{bent} configuration \((90^\circ)\) is lower in energy in this case.

b. \(\text{H}_2\text{S}\) has 8 valence electrons as well, because sulfur and oxygen are in the same periodic group. Therefore, the \textit{bent} configuration will be favored.

c. \(\text{H}_2\text{Be}\) has 4 valence electrons, which corresponds to the 2\(^{nd}\) highest orbital on the diagram. The \textit{linear} configuration \((180^\circ)\) is lower in energy in this case.

**Q10.19**

Use the Walsh diagram for the valence electrons of a \(\text{XY}_2\) molecule to predict whether the following molecules are linear or bent:
a. (CO$_2$)  b. (CO$_2^+$)  c. (CO$_2^-$)  d. (SO$_2^-$)  e. (CF$_2^+$)

### S10.19

<table>
<thead>
<tr>
<th>Valence Electrons</th>
<th>Geometry</th>
</tr>
</thead>
<tbody>
<tr>
<td>2-16</td>
<td>linear</td>
</tr>
<tr>
<td>17-20</td>
<td>bent</td>
</tr>
<tr>
<td>21-24</td>
<td>linear</td>
</tr>
</tbody>
</table>

a. CO$_2$ 16 valence electrons **Linear**  
b. CO$_2^+$ 15 valence electrons **Linear**  
c. CO$_2^-$ 17 valence electrons **Bent**  
d. SO$_2^-$ 19 valence electrons **Bent**  
e. CF$_2^+$ 18 valence electrons **Bent**
Q10.20

Walsh correlation diagrams can be used to predict the shapes of polyatomic molecules that contain more than three atoms. In this and the following three problems we consider molecules that have the general formula \(\text{XH}_3\). We will restrict our discussion to \(\text{XH}_3\) molecules, where all the \(\text{H—X—H}\) bond angles are the same. If the molecule is planar, then the \(\text{H—X—H}\) bond angle is 120°. A nonplanar \(\text{XH}_3\) molecule, then, has an \(\text{H—X—H}\) bond angle that is less than 120°. Figure 10.26 shows the Walsh correlation diagram that describes how the energies of the molecular orbitals for an \(\text{XH}_3\) molecule change as a function of the \(\text{H—X—H}\) bond angle. Note that because \(\text{XH}_3\) is not linear, the labels used to describe the orbitals on the two sides of the correlation diagram do not have designations such as

and

. We see that the lowest-energy molecular orbital is insensitive to the \(\text{H—X—H}\) bond angle. Which atomic orbital(s) contribute to the lowest-energy molecular orbital? Explain why the energy of this molecular orbital is insensitive to changes in the \(\text{H—X—H}\) bond angle.

S10.20

The lowest energy molecular orbital is the 1s orbital, which is a core atomic orbital instead of a bonding atomic orbital.

Q10.21

Consider the \(\text{BH}_2\) where the general Walsh diagram for a \(\text{XH}_2\) is shown below. What is the geometric preference of the molecule in ground and excited state?

S10.21

The \(\text{BH}_2\) molecule has the same geometric shape as water, it is bent where the HOMO is \(\psi_{\pi_u}\). The first excited state relies on the degree of bending, and the \(2a_1\) is unoccupied, and the next \(1b_2\) is the highest occupied where the preferred geometry is linear, so at the first excited state will be linear.

\(\text{BH}_2\) is a linear molecule. It has 4 valance electrons which fill 2 of the lines in the Walsh diagram. This second line has lower energy towards linear conformation and \(1\psi_{\sigma_u}\).

Q10.24

Solve for \(\psi_{\pi}\) corresponding to the energy \(E = \alpha + \beta\) for ethene.
The bonding Huckel molecular orbitals is
\[ \psi_\pi = c_1 \; 2p_{z1} + c_2 \; 2p_{z2} \]
the relationship of the coefficients can be defined as the following from the secular determinate:
\[ \begin{align*}
    &c_1 \; (\alpha - E) + c_2 \; \beta = 0 \\
    &c_1 \; \beta + c_2 \; (\alpha - E) = 0
\end{align*} \]
Substituting \( E = \alpha + \beta \) into these expressions and solving gives
\[ \begin{align*}
    &-c_1 \beta + c_2 \beta = 0 \\
    &c_1 = c_2
\end{align*} \]
Then plugging back into the original equation gives
\[ \psi_\pi = c_1 \; (2p_{zA} + 2p_{zB}) \]
Now we can solve for \( c_1 \) by normalizing the wavefunction
\[ c_1^2 (1 + 2S + 1) = 1 \]
where \( S = 0 \) so solving yields
\[ c_1 = \frac{1}{\sqrt{2}} \]
The final wave function can be written as
\[ \psi_\pi = \frac{1}{\sqrt{2}} (2p_{z1} + 2p_{z2}) \]

**Q10.25**

Generalize the molecular orbital treatment of propene allyl cation. Find the energies and wave function of this molecule.

**S10.25**

The Huckel secular determinant for propene is
\[ \begin{vmatrix}
\alpha - E & \beta & 0 \\
\beta & \alpha - E & \beta \\
0 & \beta & \alpha - E
\end{vmatrix} = 0 \]
making a substitution for \( x = \frac{\alpha - E}{\beta} \) the secular determinant becomes
\[ \begin{vmatrix}
x & 1 & 0 \\
1 & x & 1 \\
0 & 1 & x
\end{vmatrix} \]
right. = 0
Solving the determinant yields a cubic polynomial \(x^3 - 2x = 0\) the roots of this polynomial are \(x = 0, \pm \sqrt{2}\).

Replacing \(x\) with the previous substitution made it is found that \(\left[E = \alpha \pm \sqrt{2} \beta \right] \quad \text{and} \quad \left[E = 0\right]\)

To find the wave function of propene you must find the constants \(\begin{vmatrix} c_{1}(x) & c_{2} & 0 \\ c_{1} & c_{2}(x) & c_{3} \\ 0 & c_{2} & c_{3}(x) \end{vmatrix} = 0\)

solving the determinant yields \(c_{1}(1) = c_{3}(3) = \frac{1}{2}\) \(c_{2} = \frac{1}{\sqrt{2}}\)

therefore the wave function is \(\psi = \frac{1}{2} 1s + \frac{1}{\sqrt{2}} 2s + \frac{1}{2} 2p_{z}\)

Q10.26

Show that the six molecular orbitals for Benzene consturcted from the \((2Pp_{x})\) atomic orbital on each of the six carbon atoms:

\[
\psi_{i} = \sum_{j=1}^{6} c_{ij}2p_{xj}\]

leads to a secular determinant.

S10.26

The above equation for benzene is:

\[
\psi_{i} = c_{i1}2p_{z1} + c_{i2}2p_{z2} + c_{i3}2p_{z3} + c_{i4}2p_{z4} + c_{i5}2p_{z5} + c_{i6}2p_{z6}\]

The secular determinant for the benzene is:

\[
\begin{vmatrix}
H_{11} - ES_{11} & H_{12} - ES_{12} & H_{13} - ES_{13} & H_{14} - ES_{14} & H_{15} - ES_{15} & H_{16} - ES_{16} \\
H_{12} - ES_{12} & H_{22} - ES_{22} & H_{23} - ES_{23} & H_{24} - ES_{24} & H_{25} - ES_{25} & H_{26} - ES_{26} \\
H_{13} - ES_{13} & H_{23} - ES_{23} & H_{33} - ES_{33} & H_{34} - ES_{34} & H_{35} - ES_{35} & H_{36} - ES_{36} \\
H_{14} - ES_{14} & H_{24} - ES_{24} & H_{34} - ES_{34} & H_{44} - ES_{44} & H_{45} - ES_{45} & H_{46} - ES_{46} \\
H_{15} - ES_{15} & H_{25} - ES_{25} & H_{35} - ES_{35} & H_{45} - ES_{45} & H_{55} - ES_{55} & H_{56} - ES_{56} \\
H_{16} - ES_{16} & H_{26} - ES_{26} & H_{36} - ES_{36} & H_{46} - ES_{46} & H_{56} - ES_{56} & H_{66} - ES_{66}
\end{vmatrix} = 0
\]

\(H_{ij} = H_{ji}\) is a hermitian operator when \(i\) and \(j\) are neighbors

\(H_{ij} = 0\) when \(i\) and \(j\) are not neighbors

\(S_{11} = S_{22} = S_{33} = S_{44} = S_{55} = S_{66} = 1\)

\(S_{ij} = 0\)