This tutorial deals with the interpretation of the vibrational and electronic spectra of benzene using group theory.

**Analysis of the Vibrational Spectrum**

The vibrational analysis uses the method of unmoved atoms (uma). The number of unmoved atoms for each symmetry operation is stored in a vector, $\Gamma_{uma}$. This is converted into a reducible representation ($\Gamma_{vib}$) for the vibrational degrees of freedom in several steps as is shown below. Next the irreducible representations that contribute to $\Gamma_{vib}$ is determined.

\[
\begin{align*}
D_{6h} &= \begin{pmatrix} 1 & 2 & 2 & 1 & 3 & 3 & 1 & 2 & 2 & 1 & 3 & 3 \end{pmatrix} \\
D_{6h} &= D_{6h}^T & h = \sum D_{6h} & h = 24
\end{align*}
\]

\[
\begin{array}{c}
E & & E & C_6 & C_3 & C_2 & C_2' & C_2'' & i & S_3 & S_6 & \sigma_h & \sigma_d & \sigma_v \\
\Gamma_{uma} &= \begin{pmatrix} 12 & 0 & 0 & 0 & 4 & 0 & 0 & 0 & 0 & 12 & 0 & 4 \end{pmatrix}
\end{array}
\]

\[
\Gamma_{trans} = (CD_{6h}^T)^{<g>} + (CD_{6h}^T)^{<11>} & \Gamma_{rot} = (CD_{6h}^T)^{<2>} + (CD_{6h}^T)^{<5>} & \Gamma_{tot} = \overrightarrow{(\Gamma_{uma} \Gamma_{trans})}
\]

\[
\Gamma_{vib} = \Gamma_{rot} - \Gamma_{trans} - \Gamma_{rot} & i = 1 .. 12 & \text{Vib}_i = \frac{\sum \overrightarrow{D_{6h} (CD_{6h}^T)^{<i>} \Gamma_{vib}}}{h}
\]

The symmetry of the vibrational modes and their IR and Raman activity are given below:

**IR active:** $A_{2u}$ and $3E_{1u}$

**Raman active:** $2A_{1g}, E_{1g},$ and $4E_{2g}$

IR active modes are observed at 675, 1035, 1479, and 3036 cm$^{-1}$, which is consistent with the above analysis. The Raman spectrum is not as clearly resolved.
Analysis of the Electronic Spectrum

The electronic spectrum to be analyzed (see below) is due to transitions involving benzene's π electrons. The symmetry of the relevant π-electron molecular orbitals is determined by examining how the π orbitals transform under the symmetry operations of the $D_{6h}$ group.

\[
\Gamma_{\pi} = \begin{pmatrix} 6 & 0 & 0 & -2 & 0 & 0 & 0 & -6 & 0 & 2 \end{pmatrix} \quad \Pi_i = \frac{\sum \left[ D6h \ (CD6h^T)^{<i>} \tau_{\pi}^T \right]}{h}
\]

\[
\Pi = \begin{pmatrix} 0 & 0 & 0 & 1 & 1 & 0 & 0 & 1 \end{pmatrix} \quad \begin{array} \text{A1g:} x^2 + y^2 + z^2 \text{ A2g: Rz} \text{ B1g} \text{ B2g} \text{ E1g: (Rx, Ry), (xz, yz)} \text{ E2g: } (x^2 - y^2, -xy) \text{ A1u:} \text{ A2u: z} \text{ B1u:} \text{ B2u:} \text{ E1u: (x, y)} \text{ E2u:} \end{array}
\]

The symmetry of the π-molecular orbitals is $\Gamma_\pi = B_{2g} + E_{1g} + A_{2u} + E_{2u}$. The order of the levels from a Hückel calculation is as shown above: $A_{2u}$, $E_{1g}$, $E_{2u}$, and $B_{2g}$. The ground electronic state is $A_{2u}(2)$, $E_{1g}(4)$ and has $A_{1g}$ symmetry because the $A_{2u}$ and $E_{1g}$ orbitals are full. The first electronic excited state is $A_{2u}(2)$, $E_{1g}(3)$, $E_{2u}(1)$. This has the symmetry properties of $E_{1g}(1)E_{2u}(1)$ which gives rise to the manifold of states: $B_{1u}$, $B_{2u}$, and $E_{1u}$ as is shown below.

\[
\Xi_i = \frac{\sum \left[ D6h \ (CD6h^T)^{<i>} (CD6h^T)^{<5>} (CD6h^T)^{<12>} \right]}{h} \quad \Xi^T = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 & 1 & 0 \end{pmatrix}
\]

The electronic energy level diagram consistent with this analysis is shown below. Note that the excited state splits into a set of singlet and triplet exited states. Transitions from the singlet ground state to the triplet excited states are formally forbidden.
For an electronic transition to be allowed the transition moment integral must be greater than zero.

\[ \int \Psi_{ex} \mu_e \Psi_{eg} d\tau_e > 0 \]

Only the \( A_{1g} \rightarrow E_{1u} \) transition is orbitally allowed as is shown below.

\[
\begin{align*}
A_{1g} \rightarrow B_{1u} & \quad \frac{\sum \overrightarrow{\left[D6h (CD6h^T)^{<9>} \left[ (CD6h^T)^{<g>} + (CD6h^T)^{<11>} \right] (CD6h^T)^{<1>}}}{h} = 0 \\
A_{1g} \rightarrow B_{2u} & \quad \frac{\sum \overrightarrow{\left[D6h (CD6h^T)^{<10>} \left[ (CD6h^T)^{<g>} + (CD6h^T)^{<11>} \right] (CD6h^T)^{<1>}}}{h} = 0 \\
A_{1g} \rightarrow E_{1u} & \quad \frac{\sum \overrightarrow{\left[D6h (CD6h^T)^{<11>} \left[ (CD6h^T)^{<g>} + (CD6h^T)^{<11>} \right] (CD6h^T)^{<1>}}}{h} = 0
\end{align*}
\]

Electronic transitions that are orbitally forbidden can occur if they are properly coupled to vibrational transitions. This occurs when the following integral is non-zero. These are called vibronic or vibrationally assisted electronic transitions.

\[ \int \int \Psi_{ex} \Psi_{vx} \mu_e \Psi_{eg} \Psi_{vg} d\tau_e d\tau_v \]

The orbitally forbidden \( A_{1g} \rightarrow B_{1u} \) is vibronically assisted by \( B_{2g} \) or \( E_{2g} \) vibrations.

\[
X_i = \frac{\sum \overrightarrow{\left[D6h (CD6h^T)^{9} (CD6h^T)^{i} \left[ (CD6h^T)^{g} + (CD6h^T)^{11} \right] (CD6h^T)^{1} (CD6h^T)^{1}}}{h} \]

\[ X^T = \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \]

The orbitally forbidden \( A_{1g} \rightarrow B_{2u} \) is vibronically assisted by \( B_{1g} \) or \( E_{2g} \) vibrations. However there is no \( B_{1g} \) vibration.

\[
X_i = \frac{\sum \overrightarrow{\left[D6h (CD6h^T)^{10} (CD6h^T)^{i} \left[ (CD6h^T)^{g} + (CD6h^T)^{11} \right] (CD6h^T)^{1} (CD6h^T)^{1}}}{h} \]

\[ X^T = \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \]
The fully allowed $A_{1g} \rightarrow E_{1u}$ transition is assigned to the most intense transition which occurs at 180 nm. The vibronically assisted $A_{1g} \rightarrow B_{1u}$ and $A_{1g} \rightarrow B_{2u}$ transitions are assigned to the less intense bands at 200 and 260 nm, respectively. The spin-forbidden $1 A_{1g} \rightarrow 3 B_{1u}$ is assigned to the lowest energy and lowest intensity transition at 340 nm.